

WRF

Coriolis Force in the WRF model

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Equations

- ◆ $U(t+\Delta t) = U(t) + \Delta t * F_{U,Cor}$
- ◆ $V(t+\Delta t) = V(t) + \Delta t * F_{V,Cor}$
- ◆ $W(t+\Delta t) = W(t) + \Delta t * F_{W,Cor}$

- ◆ This is a forward in time method
- ◆ The WRF model describes both the horizontal and vertical Coriolis forces

Force Equations

- ◆ The forces $F_{U,Cor}$, $F_{V,Cor}$, $F_{W,Cor}$ are described by the following equations:
- ◆
$$F_{U,Cor} = [(f_{i+1/2} + f_{i-1/2})/2] * [(V_{i+1/2,j+1/2} + V_{i+1/2,j-1/2} + V_{i-1/2,j+1/2} + V_{i-1/2,j-1/2})/4] - [(e_{i+1/2} + e_{i-1/2})/2] * [W_{i+1/2,k+1/2} + W_{i+1/2,k-1/2} + W_{i-1/2,k+1/2} + W_{i-1/2,k-1/2})/4] * [(cos\alpha_{i+1/2} + cos\alpha_{i-1/2})/2]$$
- ◆
$$F_{V,Cor} = -[(f_{j+1/2} + f_{j-1/2})/2] * [(U_{i+1/2,j+1/2} + U_{i+1/2,j-1/2} + U_{i-1/2,j+1/2} + U_{i-1/2,j-1/2})/4] - [(e_{j+1/2} + e_{j-1/2})/2] * [W_{j+1/2,k+1/2} + W_{j+1/2,k-1/2} + W_{j-1/2,k+1/2} + W_{j-1/2,k-1/2})/4] * [(sin\alpha_{j+1/2} + sin\alpha_{j-1/2})/2]$$
- ◆
$$F_{W,Cor} = e * \{ [(U_{i+1/2,k+1/2} + U_{i+1/2,k-1/2} + U_{i-1/2,k+1/2} + U_{i-1/2,k-1/2})/4] * cos\alpha - [(V_{j+1/2,k+1/2} + V_{j+1/2,k-1/2} + V_{j-1/2,k+1/2} + V_{j-1/2,k-1/2})/4] * sin\alpha \}$$

Force Equations

- ◆ These equations reduce to the following when we remove the grid staggering:
- ◆ $F_{U,Cor} = f*V - e*W*cos\alpha$
- ◆ $F_{V,Cor} = -f*U - e*W*sin\alpha$
- ◆ $F_{W,Cor} = e*U*cos\alpha - V*sin\alpha$

Force Equations

- ◆ This leaves us with the equations for the Coriolis Force as:
- ◆ $U(t+\Delta t) = U(t) + \Delta t * (f * V - e * W * \cos \alpha)$
- ◆ $V(t+\Delta t) = V(t) + \Delta t * (-f * U - e * W * \sin \alpha)$
- ◆ $W(t+\Delta t) = W(t) + \Delta t * (e * U * \cos \alpha - V * \sin \alpha)$

Force Equations

- ◆ where α is the local rotation angle between the y -axis and the meridians
- ◆ $e = 2\Omega \cos\varphi$
- ◆ $f = 2\Omega \sin\varphi$
- ◆ φ is the latitude
- ◆ Includes both horizontal and vertical effects

Stability Analysis

- ◆ Let $u = u_0 \exp[i(kn\Delta x + \omega\tau\Delta t)]$
- ◆ $v = v_0 \exp[i(kn\Delta x + \omega\tau\Delta t)]$
- ◆ $w = w_0 \exp[i(kn\Delta x + \omega\tau\Delta t)]$
- ◆ and $\Psi = \exp(i\omega\Delta t)$. Plugging these values into the equations and putting it in matrix form we get:

Stability Analysis

$$\begin{vmatrix} \Psi-1 & -f \Delta t & e \cos \alpha \Delta t \\ f \Delta t & \Psi-1 & -e \sin \alpha \Delta t \\ -e \cos \alpha \Delta t & e \sin \alpha \Delta t & \Psi-1 \end{vmatrix} \begin{vmatrix} u_0 \\ v_0 \\ w_0 \end{vmatrix} = \begin{vmatrix} 0 \\ 0 \\ 0 \end{vmatrix}$$

- ◆ We need to find the determinant of the matrix and set it equal to zero

Stability Analysis

- ◆ When we set the determinant equal to zero we get

$$(\Psi-1)[(\Psi-1)^2 + e^2\Delta t^2 + f^2\Delta t^2] = 0$$

- ◆ Now we need to solve for Ψ

Stability Analysis

- ◆ Solving for we get $\Psi = 1, 1 \pm i \Delta t \sqrt{(e^2+f^2)}$
- ◆ Equating real and imaginary parts we get:
$$\lambda \cos(\omega_r \Delta t) = 1$$
$$\lambda \sin(\omega_r \Delta t) = \pm \Delta t \sqrt{(e^2+f^2)}$$
- ◆ Squaring and summing we find that
$$\lambda^2 = 1 + \Delta t^2(e^2+f^2)$$

Stability Analysis

- ◆ Solving for λ we find that
$$\lambda = \sqrt{[1 + \Delta t^2(e^2 + f^2)]} \geq 1$$
- ◆ This shows that $\lambda = 1$ only when $\Delta t = 0$
- ◆ The scheme for the Coriolis force is unstable in WRF

References

- Pielke, R.A., Sr., 2002: Mesoscale meteorological modeling. 2nd Edition, Academic Press, San Diego, CA, 676 pp.
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